

Vibrational Transition Moments and Dipole Derivatives

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Oscillator strengths and integrated absorption intensities of simple vibrational absorption (infrared) spectra are related to the squares of electric-dipole transition moments,¹

$$D_{vv'}^n = \langle \Psi_{nv} | \hat{\mu} | \Psi_{nv'} \rangle \cdot \langle \Psi_{nv'} | \hat{\mu} | \Psi_{nv} \rangle, \quad (1)$$

where n denotes the electronic state, v and v' denote vibrational states, and $|\Psi_{nv}\rangle$ and $|\Psi_{nv'}\rangle$ denote initial and final vibronic states, respectively.

We may compute the electric-dipole vibrational transition moment beginning from the Born-Oppenheimer approximation, in which we assume that the total vibronic wave function, $\Psi_{nv}(\mathbf{r}, \mathbf{R})$, may be written as a product of an electronic wave function, $\psi_n(\mathbf{r}; \mathbf{R})$ and a vibrational wave function, $\chi_{nv}(\mathbf{R})$, where \mathbf{r} and \mathbf{R} denote the collective electronic and nuclear coordinates, respectively. Then the electric-dipole transition matrix element may be written as

$$\langle \Psi_{nv}(\mathbf{r}, \mathbf{R}) | \hat{\mu} | \Psi_{nv'}(\mathbf{r}, \mathbf{R}) \rangle = \langle \chi_{nv}(\mathbf{R}) | \langle \psi_n(\mathbf{r}; \mathbf{R}) | \hat{\mu} | \psi_n(\mathbf{r}; \mathbf{R}) \rangle | \chi_{nv'}(\mathbf{R}) \rangle = \langle \chi_{nv} | \langle \hat{\mu} \rangle_n | \chi_{nv'} \rangle \quad (2)$$

where $\langle \hat{\mu} \rangle_n$ denotes the expectation value of the electric-dipole operator in the n -th Born-Oppenheimer electronic state. The dependence of the $\langle \hat{\mu} \rangle_n$ on the nuclear coordinates is usually approximated by the first term of its Taylor expansion about a reference geometry \mathbf{R}^0 (*i.e.*, the *electrical harmonic approximation*):

$$\langle \hat{\mu} \rangle_n \approx \langle \hat{\mu} \rangle_0 + \sum_{\alpha} \left(\frac{\partial \langle \hat{\mu} \rangle_n}{\partial R_{\alpha}} \right)_0 (R_{\alpha} - R_{\alpha}^0), \quad (3)$$

where the subscript 0 indicates that the given quantity is evaluated at the reference geometry. The dipole-moment derivatives may be easily computed using analytic gradient techniques, and the final expressions vary depending on the level of theory employed.² The total electric-dipole transition matrix element then becomes

$$\langle \Psi_{nv} | \hat{\mu} | \Psi_{nv'} \rangle = \sum_{\alpha} \left(\frac{\partial \langle \hat{\mu} \rangle_n}{\partial R_{\alpha}} \right)_0 \langle \chi_{nv} | (R_{\alpha} - R_{\alpha}^0) | \chi_{nv'} \rangle. \quad (4)$$

The vibrational wave functions, χ_{nv} , are usually taken to be harmonic oscillator functions (*i.e.*, the *mechanical harmonic approximation*), which subsequently leads to relatively simple programmable equations in terms of the normal vibrational modes.³

References

- [1] L. D. Barron, *Molecular Light Scattering and Optical Activity*, 2nd edition ed. (Cambridge University Press, Cambridge, U.K., 2004).
- [2] Y. Yamaguchi, Y. Osamura, J. D. Goddard, and H. F. Schaefer, *A New Dimension to Quantum Chemistry: Analytic Derivative Methods in Ab Initio Molecular Electronic Structure Theory*, No. 29 in *International Series of Monographs on Chemistry* (Oxford Univ. Press, New York, 1994).
- [3] E. B. Wilson, J. C. Decius, and P. C. Cross, *Molecular Vibrations: The Theory of Infrared and Raman Vibrational Spectra* (Dover, New York, 1980).